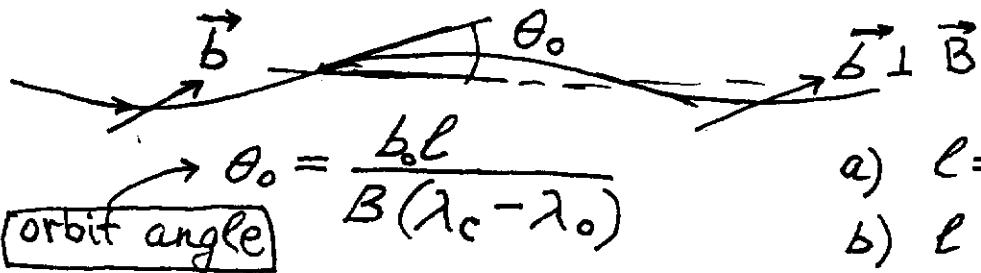


Ionization Cooling on Spiral Orbit in Solenoid

- Field composition:
 - continuous solenoid (or long solenoids) plus
 - low, quasi-resonance dipole field
- Beam track:
 - quasi-resonance spiral not exceeding beam size
 - absorbers with rotating gradient
- Results in 3-dim. cooling
- Dipole field options:
 - a) continuous helix period λ_0
 - b) one-directed lumped dipoles field, same period

$\vec{B} \rightarrow$ $\lambda_c - \lambda_0 \ll \lambda_0$ $\lambda_c = 2\pi \frac{P C}{e B}$



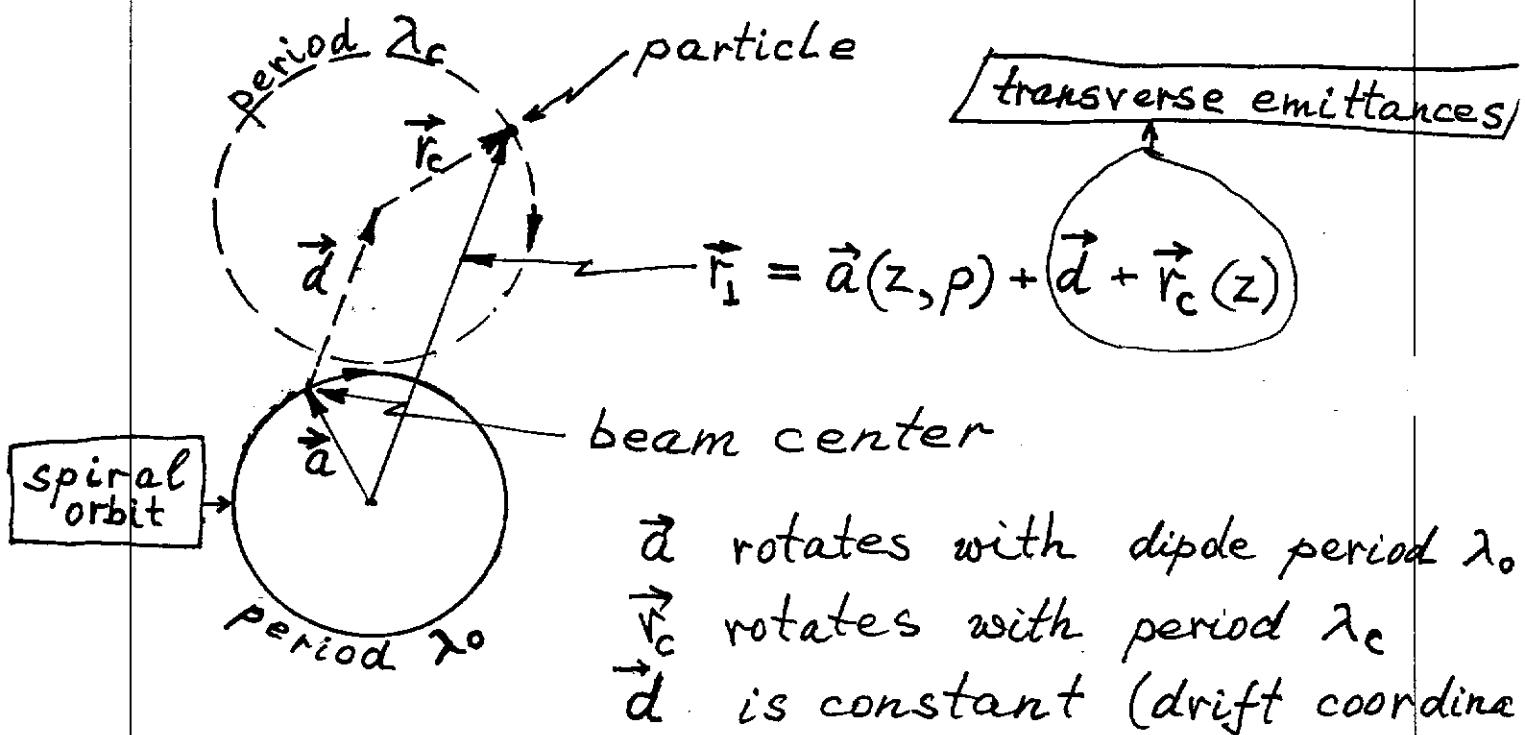
a) $l = \lambda_0$
b) $l \ll \lambda_0$

Orbit radius: $a = \theta_0 \lambda_0 / 2\pi$

Dispersion: $\vec{D} = P \frac{d\vec{a}}{dp} = -\frac{\lambda_c}{\lambda_c - \lambda_0} \vec{a}$

$$|\lambda_c - \lambda_0 \gg \lambda_0 (\theta_0^2 + \theta_c^2 + \frac{\Delta p}{P})|$$

Transverse motion:



- \vec{a} rotates with dipole period λ_0
- \vec{r}_c rotates with period λ_c
- d is constant (drift coordinate)

Absorber gradient: $\vec{\nabla}n = \frac{\vec{a}}{a} \nabla n$
(rotates)

$$\nabla n = \text{const}$$

$$n = n_0 + (\nabla n) \frac{\vec{a}}{a} \cdot \left(\vec{D} \frac{\Delta p}{P} + \vec{d} + \vec{r}_c \right)$$

Friction force:

$$\vec{f} = -E' \frac{\vec{p}}{P}; -E' = \frac{4\pi n Z e^4 \log}{m_e c^2 \beta^2}$$

($\log \approx 15$)

- Averaged friction effects:

$$\vec{r}' = -\frac{1}{2} \lambda_c \vec{r}_c ;$$

$$\vec{d}' = -\frac{1}{2} \lambda_d \vec{d} ; \quad (\text{beam shrinks!})$$

$$\Delta p' = -\lambda_{||} \Delta p .$$

Cooling decrements:

$$\lambda_{||} = \frac{|\vec{r}_0'|}{\gamma \beta^2} \left(-\frac{2}{\gamma^2} + \frac{2}{\log} - \frac{\lambda_c}{\lambda_c - \lambda_0} \frac{\alpha \nabla n}{n_0} \right)$$

$$\lambda_d = -\frac{|\vec{r}_0'|}{\gamma \beta^2} \frac{\lambda_c}{\lambda_0} \frac{\alpha \nabla n}{n_0}$$

$$\lambda_c = \frac{|\vec{r}_0'|}{\gamma \beta^2} \left[2 + \frac{\lambda_c^2}{\lambda_0(\lambda_c - \lambda_0)} \frac{\alpha \nabla n}{n_0} \right]$$

- Decrement sum:

$$\lambda_6 = \lambda_c + \lambda_{||} + \lambda_d = 2 \frac{|\vec{r}_0'|}{\gamma} \left(1 + \frac{1}{\beta^2 \log} \right),$$

as it must be ($\lambda_6 = -\frac{\partial}{\partial P} \vec{f}(\vec{P}, \vec{r})$)

Second order effects (negligible...)

- drift tube (slow rotation): $\frac{\lambda_d}{\lambda_c} \sim \frac{1}{\theta_0^2} \frac{\lambda_c}{\lambda_c - \lambda_0}$

- decrements shifts:

$$\frac{\Delta \lambda}{\lambda} \sim \frac{\lambda_c}{\lambda_d}$$

Decrments redistribution (options)

1) Conditions for $\lambda_c = \lambda_{II} = \lambda_d = \frac{1}{3} \lambda_0 = \frac{2}{3} \beta^2$ ($\beta^2 \gg 1/\log$)

$$\left. \begin{aligned} \frac{\alpha \nabla n}{n_0} &= -\frac{2}{3} \beta^2 \frac{3-2\beta^2}{3-\beta^2} \\ \frac{\lambda_c - \lambda_0}{\lambda_c} &= \frac{\beta^2}{3-\beta^2} \end{aligned} \right\}$$

2) Equalizing between λ_c and λ_{II} :

$$\lambda_c = \lambda_{II} = \frac{|\gamma'_0|}{\gamma}, \quad \lambda_d = -\lambda_c \frac{\alpha \nabla n}{n_0}$$

at

$$0 < -\frac{\alpha \nabla n}{n_0} = 2 \frac{\lambda_0}{\lambda_c} \frac{\lambda_c - \lambda_0}{\lambda_c + \lambda_0} (2 - \beta^2) \ll \beta^2$$

3) Total emittance exchange:

$$\lambda_c \ll \lambda_{II}, \quad \lambda_d \ll \lambda_{II}$$

$$\lambda_{II} \approx 2 \frac{|\gamma'_0|}{\gamma}$$

at

$$0 < -\frac{\alpha \nabla n}{n_0} = 2 \frac{\lambda_c - \lambda_0}{\lambda_c^2} \lambda_0 \ll 2\beta^2$$

Important outlines:

- parameter $\frac{\alpha \nabla n}{n_0}$ can always be small
- hence, $\theta_0 = 2\pi a / \lambda_0$ can be quite small, as well.

Dipole field estimation

assume $a = 6 \Rightarrow \sqrt{\epsilon \cdot mc^2/eB}$

then

$$\frac{b_0}{B} = \frac{\lambda_c - \lambda_0}{\lambda_0} \cdot 2\pi \frac{c}{l} \approx \frac{\lambda_c - \lambda_0}{\lambda_0} \cdot \frac{\lambda_0}{l} \theta_c$$

$$\left. \begin{array}{l} \epsilon = 1.5 \text{ cm} \\ B = 6 \text{ T} \\ \gamma = 2 \end{array} \right\} \rightarrow \theta_c = \frac{1}{\gamma \beta} \sqrt{\epsilon \cdot \frac{eB}{mc^2}} = 0.25$$
$$\lambda_c = 12 \text{ cm} \times 2\pi \approx 75 \text{ cm}$$

assume $\frac{\lambda_c - \lambda_0}{\lambda_0} = 0.25$ ("quasi-resonance")

a) $l = \lambda_0$ (continuous helix)

then $\frac{b_0}{B} = 0.0625$; $b_0 \approx 4 \text{ kG}$

b) $l = \lambda_0/4$, then $b_0 = 1.6 \text{ T}$

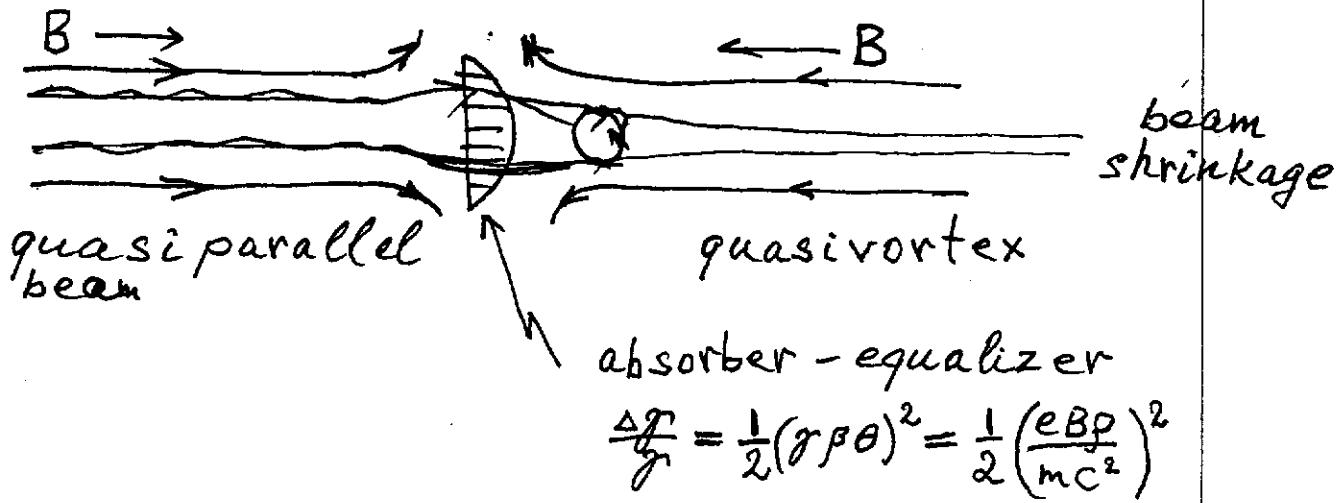
- Required $b_0 l$ value reduces with cooling

Options for NF

- Mini - Cooling with total emittance exchange (i.e. $\lambda_c \approx 0$, $\lambda_d \approx 0$):

$$\Delta p = \Delta p_{in} \cdot (\gamma/\gamma_{in})^2$$

- Second cooling stage: $\lambda_{\parallel} = \lambda_c$; $\lambda_d \approx 0$
two long solenoids



- Final stage: $\lambda_{\parallel} = \lambda_c = \lambda_d$
(or different optimization, to obtain minimum E_6)

- Post-equilibrium stage:
reverse fast emittance exchange
- $$\lambda_{\parallel} < 0, \quad |\lambda_{\parallel}| \gg \lambda_s$$
- $$\lambda_c + \lambda_d \gg \lambda_s$$

Conclusions

- Spiral cooling concept seems efficient and flexible:
 - Continuous solenoid
 - Non-difficult dipole arrangement
 - Stable beam
 - Effective emittance exchange
 - 3-dimensional cooling
 - Mini-Cooling at total emittance exchange seems very important possibility — to make it easier the beam RF capture in both projects, Neutrino Factory and Muon Collider
- Spiral Cooling seems capable to provide the minimum 6-dim emittance (reduction by a factor $\approx 10^6$).